## <sup>1</sup>H NMR Studies on the Molecular Dynamics of Acetylferrocene in Crystals

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Z. Naturforsch. 43a, 78-80 (1988); received August 14, 1987

The temperature dependence of  $^1\mathrm{H}$  spin-lattice relaxation time was determined at 20 MHz for solid acetylferrocene [( $\mathrm{C_5H_5}$ )( $\mathrm{C_5H_4COCH_3}$ )Fe] from ca. 80 K up to the m.p. (359 K). Rather large activation energies of 21 and 24 kJ mol  $^{-1}$  for the  $\mathrm{C_5}$  reorientations of the two crystallographically nonequivalent non-substituted cyclopentadienyl rings were obtained, indicating that the crystal has a closely packed structure. The two kinds of  $\mathrm{CH_3}$  groups attached to the substituted cyclopentadienyl rings were assumed to be approximately equivalent and gave a low activation energy of 4 kJ mol  $^{-1}$  for the  $\mathrm{C_3}$  reorientation. No phase transition was observed in the relaxation times or in additional experiments of differential thermal analysis although the presence of two phase transitions has been reported previously at temperatures immediately below the melting temperature by means of differential scanning calorimetry.

## Introduction

Recently we have carried out  $^1H$  NMR studies on various carbocyclic  $\pi$  complexes, i.e., ferrocene, azaferrocene, ruthenocene, etc., to clarify phase transitions, the crystal structure of newly found phases, and the dynamics of the molecule as a whole and the five membered rings [1-4]. We found that in the crystals the cyclopentadienyl  $(C_5H_5; C_p)$  and pyridinyl  $(C_4H_4N)$  rings perform  $C_5$  and pseudo- $C_5$  reorientations, respectively, with the quite different activation energies being mostly determined by intermolecular interactions. This means that, as for the low-temperature stable and substable ferrocene crystals [4], crystal packing is the major origin for the barrier to rotation of the rings.

Acetylferrocene, having an acetyl group attached to one of the  $C_{\rm p}$  rings, is expected to form a rather loosely packed crystal although the bulky acetyl group may create a hindering barrier for the rotation of the  $C_{\rm p}$  rings. Sato et al. [5] studied the crystal structure of this compound at room temperature and below the melting point and proposed that the molecules rotate around the central iron atom at higher temperatures. Early  $^1{\rm H}$  NMR investigations, however, had indicated reorientation of the non-substituted  $C_{\rm p}$  rings rather than that of the molecules [5 a]. The present investigation has been undertaken to clarify the molecular dynamics of the  $\pi$  complex of solid acetylferrocene.

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## **Experimental**

Commercial acetylferrocene was purified three times by sublimation at ca. 60° C under reduced pressure. X-ray powder patterns by means of an X-ray diffractometer Model D-3F from Rigaku Denki Co. showed that the crystals of the present investigation have exactly the same structure as those employed in the experiments of Sato et al. [5].

The spin-lattice relaxation time  $T_1$  of the <sup>1</sup>H nuclei was determined at various temperatures between 80 and 359 K (m.p.) by use of a homemade pulsed NMR spectrometer [6] operated at the Larmor frequency 20 MHz. A pulse sequence of  $180^{\circ}$ – $\tau$ – $90^{\circ}$  was employed for the determination of  $T_1$ . The  $T_1$  values observed were estimated to be accurate within  $\pm$  5% and the temperatures within  $\pm$  1 K. The differential thermal analysis (DTA) was carried out by use of a homemade apparatus already described [7].

## **Results and Discussion**

According to Sato et al. [5], acetylferrocene forms monoclinic crystals belonging to the space group  $P2_1/c$  with Z=8.

From the experiments of differential scanning calorimetry (DSC), Sato et al. [5] found with a very slow heating rate of 0.017 K min<sup>-1</sup> two successive phase transitions immediately below the melting point. In our DTA experiments, carried out repeatedly with a heating rate of 0.07 K min<sup>-1</sup>, however, no heat anom-

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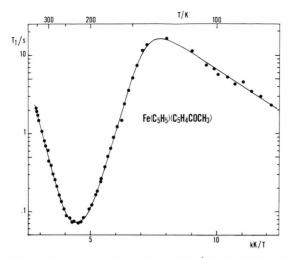


Fig. 1. Temperature dependence of the  $^1\mathrm{H}$  spin-lattice relaxation time  $T_1$  observed at 20 MHz for acetylferrocene. Solid line: calculated  $T_1$  curve.

aly except that at the m.p. was found. The melting temperature 358 K agreed very well with that reported by Sato et al. [5].

Figure 1 shows the obtained temperature dependence of  ${}^{1}H$   $T_{1}$  between 80 K and the m.p. at the Larmor frequency of 20 MHz. The  $T_{1}$  curve shows no anomaly just below the m.p., confirming our DTA result, it shows a minimum of 0.07 s at ca. 220 K, a slightly distorted V-shaped curve as compared to the usual BPP one [8], and a maximum of 20 s around 130 K. Below this maximum  $\log T_{1}$  decreases almost lineary with  $T^{-1}$ .

The  $^1$ H  $T_1$  curve indicates that at least two different kinds of motional processes successively occur with increasing temperature and that they partly overlap to form a  $T_1$  maximum. The gentle  $\log T_1$  decrease with increasing  $T^{-1}$  is most likely attributable to the CH $_3$  C $_3$  reorientations. This is because the CH $_3$  groups in the acetyl groups always have very small activation energies for the CH $_3$  C $_3$  reorientations [9, 10]. The present assignment is also consistent with the previous  $^1$ H wide-line NMR results reported by Makova et al. [11].

According to the crystal structure of acetylferrocene, two crystallographically nonequivalent molecules exist in the crystal [5]. This means that there are two kinds of  $C_p$  rings having different surroundings in the crystal. In fact, the  $T_1$  minimum attributable to the  $C_5$  reorientation of the  $C_p$  rings is slightly flattened.

Table 1. <sup>1</sup>H NMR second moment reduction  $\Delta M_2$  for the activation of the three considered modes, their activation energies  $E_a$ , and correlation times  $\tau_0$ .

| Motional Mode (i)  | $\Delta M_2^{(i)}/{ m G}^2$                   | $E_{\rm a}/{\rm kJmol^{-1}}$   | $\log (\tau_0/s)$  |
|--|---|--|--|
| $ \begin{array}{c} CH_3 \text{ reorient. (1)} \\ C_5H_5 \text{ reorient. (2)} \\ (3) \end{array} $ | 5.2 (fixed)<br>1.5 $\pm$ 0.6<br>1.6 $\pm$ 0.6 | $\begin{array}{c} 4.0 \pm 0.2 \\ 21.1 \pm 0.6 \\ 23.5 \pm 1.3 \end{array}$ | $\begin{array}{c} -13.0 \pm 0.2 \\ -13.5 \pm 0.2 \\ -13.5 \pm 0.2 \end{array}$ |

We therefore assume two sets of motional parameters for the motion of the  $C_{\rm p}$  rings, namely i=1,2. The motional parameters of the CH<sub>3</sub> C<sub>3</sub> reorientation were roughly approximated by only one set (i=3), although there must be two different kinds of CH<sub>3</sub> groups in the crystal.

Therefore, to obtain the motional parameters for the CH<sub>3</sub> C<sub>3</sub> and  $C_p$  C<sub>5</sub> reorientations, we assume that the observed <sup>1</sup>H  $T_1$  can be expressed as the sum of the three independent BPP curves due to the CH<sub>3</sub> C<sub>3</sub> and the two kinds of the  $C_p$  C<sub>5</sub> reorientations given by [12]

$$T_1^{-1} = \frac{3}{2} \gamma \sum_{i=1}^{3} \Delta M_2^{(i)} \left\{ \frac{\tau_i}{1 + \tau_i^2 \,\omega_0^2} + \frac{4 \,\tau_i}{1 + 4 \,\tau_i^2 \,\omega_0^2} \right\}, \quad (1)$$

where the correlation times  $\tau_i$  for the *i*-th motional mode are assumed to obey the Arrhenius relationships

$$\tau_i = \tau_0^i \exp\left(E_a^i / RT\right). \tag{2}$$

In (1) and (2),  $\gamma$  and  $\omega_0$  represent the proton gyromagnetic ratio and the angular Larmor frequency, respectively, and  $\Delta M_2^{(i)}$ ,  $\tau_0^i$ , and  $E_a^i$  denote the second moment reduction due to the activation of the *i*-th motional mode of a group, the correlation time of the *i*-th motional mode in the limit of infinite temperature, and its activation energy, respectively.

 $\Delta M_2^{(3)}$  was calculated to be 5.2  $G^2$  by use of the known crystal data [5]. The intermolecular contribution to  $\Delta M_2^{(3)}$  was summed up in the same way as employed in [4]. The fitting of (1) and (2) to the observed  $T_1$  data was performed by a least-squares method. The calculated curve explains the measurements very well. The motional parameters evaluated are listed in Table 1.

To check the significance of the motional parameters we calculated  $\Delta M_2^{(i)}$  for the  $C_p$   $C_5$  reorientation using the crystal data [5]. The intermolecular contribution was estimated in a similar way as that used for the evaluation of  $M_2$  of the CH<sub>3</sub> group. The  $M_2$  values calculated for the model in which one kind or two kinds of  $C_p$  rings perform  $C_5$  reorientation in addition

to the CH<sub>3</sub> C<sub>3</sub> reorientation were found to be 6.3 and 4.8 G<sup>2</sup>, respectively. This leads to  $\Delta M_2^{(i)}$  equal to 1.5 and 1.6 G<sup>2</sup>, where i is 1 or 2. So the calculated  $\Delta M_2^{(i)}$ values agree very well with the values obtained from the fitting calculation.

The observed  $E_a$  value for the  $CH_3$   $C_3$  reorientation of acetylferrocene is 4.0 kJ mol<sup>-1</sup>. This is comparable with the data obtained for acetone [10], toluene [13],  $CH_3MCl_3$  (m = Si, Ge, Sn) [14], and dimethylpyri-

- dines [15]. However,  $E_a$  values for the CH<sub>3</sub> C<sub>3</sub> reorientation of CH<sub>3</sub>NH $_3^+$  cations in various compounds studied previously [16-20] are much larger than the present value. The  $E_a$  values of  $21-23 \text{ kJ mol}^{-1}$  for the  $C_p$   $C_5$  reorientation are fairly large for this motion. They are comparable with that of the low-temperature stable orthorhombic phase of ferrocene and are larger than those of triclinic and monoclinic ferrocene.
- [1] A. Kubo, R. Ikeda, and D. Nakamura, Chem. Lett. **1981,** 1497.
- [2] A. Kubo, R. Ikeda, and D. Nakamura, Chem. Lett. 1982, 1487.
- [3] A. Kubo, R. Ikeda, and D. Nakamura, Ber. Bunsenges. Phys. Chem. **90**, 479 (1986).
- [4] A. Kubo, R. Ikeda, and D. Nakamura, J. Chem. Soc. Faraday Trans. 2, 82, 1543 (1986).
- [5] K. Sato, M. Katada, H. Sano, and M. Konno, Bull. Chem. Soc. Japan 57, 2361 (1984).
- [5a] C. H. Holm and J. A. Ibers, J. Chem. Phys. 30, 885 (1959). – L. N. Mulay and A. Attalla, J. Amer. Chem. Soc. **85**, 702 (1963).
- [6] L. S. Prabhumirashi, R. Ikeda, and D. Nakamura, Ber. Bunsenges. Phys. Chem. 85, 1142 (1981).
- [7] Y. Kume, R. Ikeda, and D. Nakamura, J. Magn. Reson. 33, 331 (1979).
- [8] N. Bloembergen, E. M. Purcell, and R. V. Pound, Phys. Rev. 73, 679 (1948).
- [9] P. S. Allen and P. Branson, J. Phys. C 11, L121 (1978).
   [10] P. S. Allen, P. Branson, M. Punkkinen, and D. G. Tayler, J. Phys. C 9, 4453 (1976).

- [11] M. K. Makova, E. V. Leonova, Yu. S. Karimov, and N. S. Kochetkowa, J. Organometal. Chem. 55, 185 (1973).
- [12] cf. e.g. A. Abragam, The Principles of Nuclear Magnetism, Oxford 1961.
- [13] A. M. I. Ahmed and R. G. Eades, Faraday Trans. II, 68, 1337 (1972).
- [14] I. Y. Wei and C. S. Johnson, Jr., J. Magn. Reson. 28, 377 (1977).
- [15] W. Müller-Warmuth, R. Schüler, M. Prager, and A. Kollmar, J. Chem. Phys. 69, 2382 (1978).
- [16] R. Ikeda, Y. Kume, D. Nakamura, Y. Furukawa, and H. Kiriyama, J. Magn. Reson. 24, 9 (1976).
- [17] H. Ishida, R. Ikeda, and D. Nakamura, J. Phys. Chem. 86, 1003 (1982).
- [18] H. Ishida, R. Ikeda, and D. Nakamura, Bull. Chem. Soc. Japan 55, 3116 (1982).
- [19] H. Ishida, R. Ikeda, and D. Nakamura, Bull. Chem. Soc. Japan 59, 915 (1986).
- [20] H. Ishida, R. Ikeda, and D. Nakamura, Bull. Chem. Soc. Japan 60, 467 (1987).